



METAXIAL DIFFRACTION METHOD

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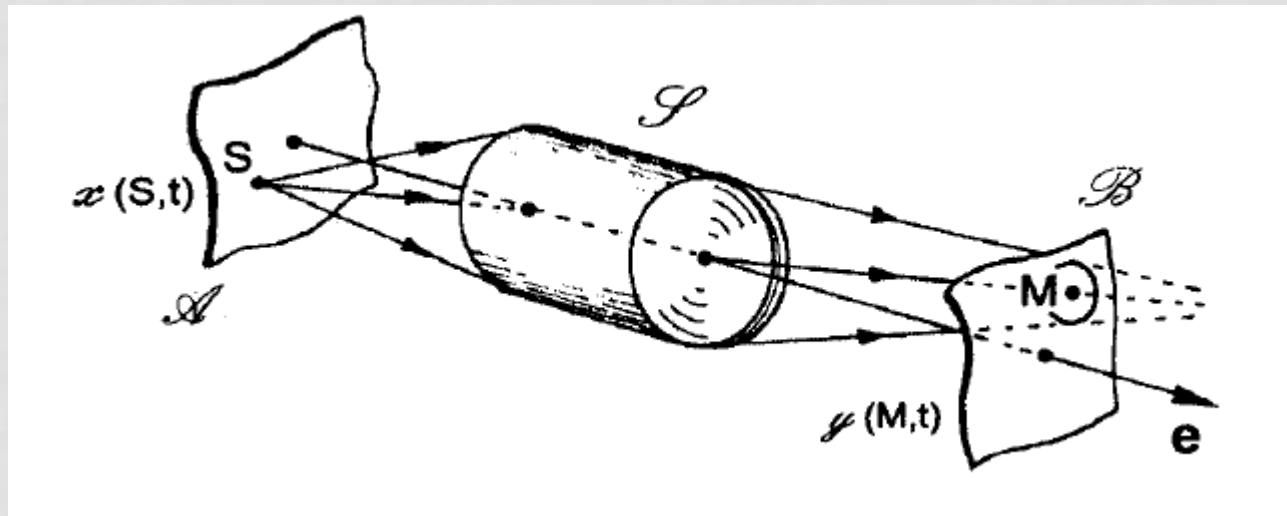
The Need of Unified Framework

- Optical communications is a must in broadband era.
- Optical communications can be free space or through optical fiber
- Fundamental phenomena: field transfer from transmitter to the receiver through electromagnetic system
- In this presentation, we propose unified framework to analyse and design of optical systems – Metaxial Optics-- developed by G. Bonnet and P. Pellat-Finet.

SPATIAL DOMAIN FIELD TRANSFER

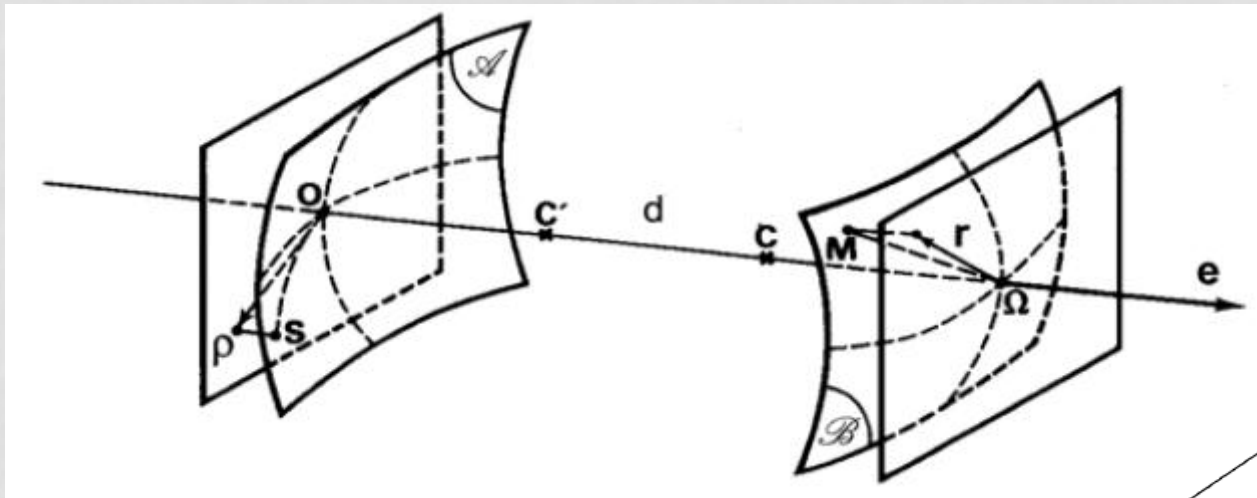
PRINCIPLES

- Radiative transfer between transmitter and receiver traversing an electromagnetic system
- Geometric : spherical transmitter, receiver, and wavefront
 - plane is a special case
- Field transfer system

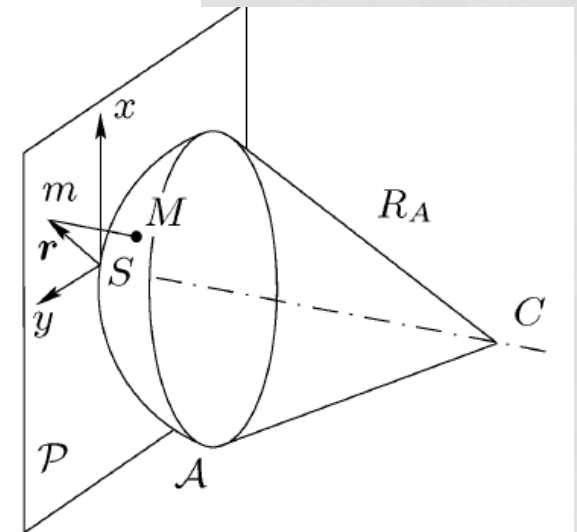


GENERAL SPATIAL PROBLEM

- Field transfer from spherical A to spherical B

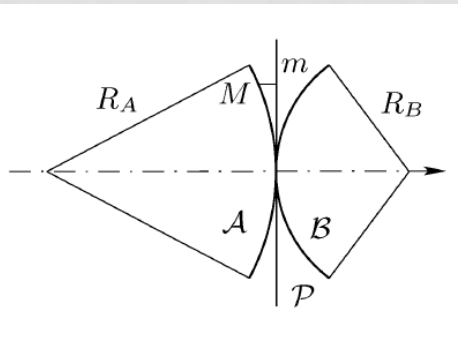


- Transmitter coordinate



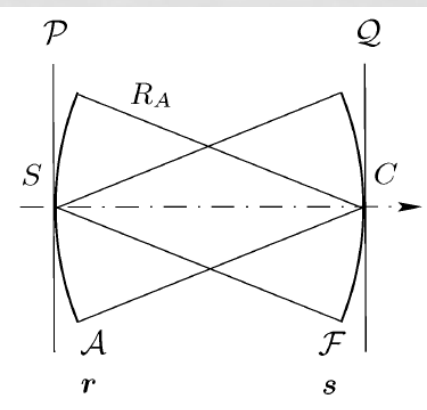
PHASE AND FOURIER TRANSFORMATIONS

- Phase transformation (curvature transparency)
 - Field transfer between two tangential sphere



$$U_B(\mathbf{r}) = U_A(\mathbf{r}) \exp \left[-\frac{i\pi}{\lambda} \left(\frac{1}{R_B} - \frac{1}{R_A} \right) r^2 \right]$$

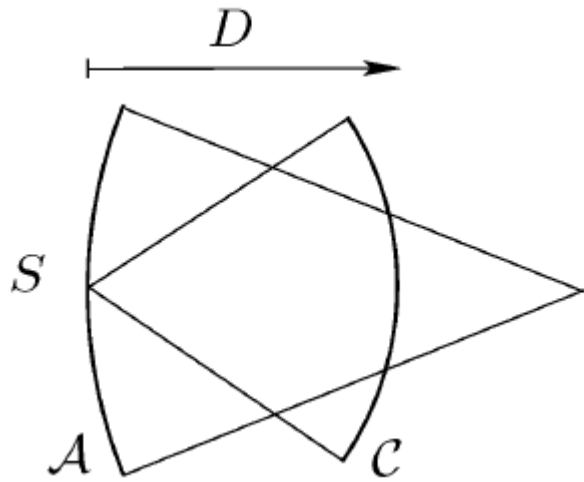
- Fourier transformation:
 - field transfer between two conjugate sphere



$$U_F(\mathbf{s}) = \frac{i}{\lambda R_A} \hat{U}_A \left(\frac{\mathbf{s}}{\lambda R_A} \right) = \frac{i}{\lambda D} \hat{U}_A \left(\frac{\mathbf{s}}{\lambda D} \right)$$

We omit propagation delay factor $\exp \left[-\frac{2i\pi R_A}{\lambda} \right]$

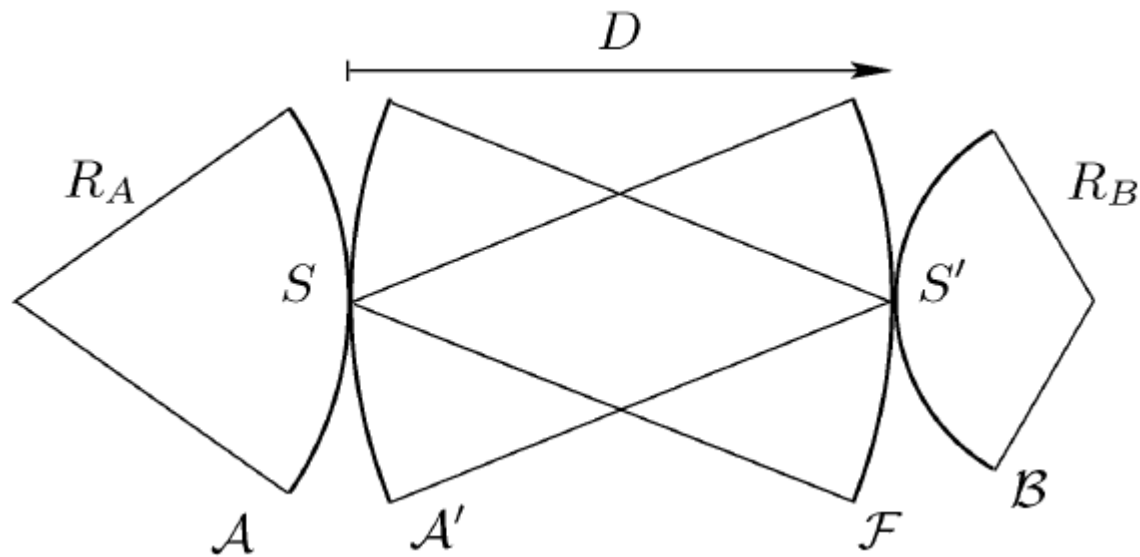
SPECIAL CASE: FRESNEL DIFFRACTION



Field transfer between cardinal sphere :
From sphere A to sphere C

$$U_C(\mathbf{s}) = \frac{i}{\lambda D} \int_{\mathbb{R}^2} \exp \left[-\frac{i\pi}{\lambda} \left(\frac{1}{D} - \frac{1}{R_A} \right) r^2 \right] \exp \left[\frac{2i\pi}{\lambda D} \mathbf{s} \cdot \mathbf{r} \right] U_A(\mathbf{r}) d\mathbf{r}$$

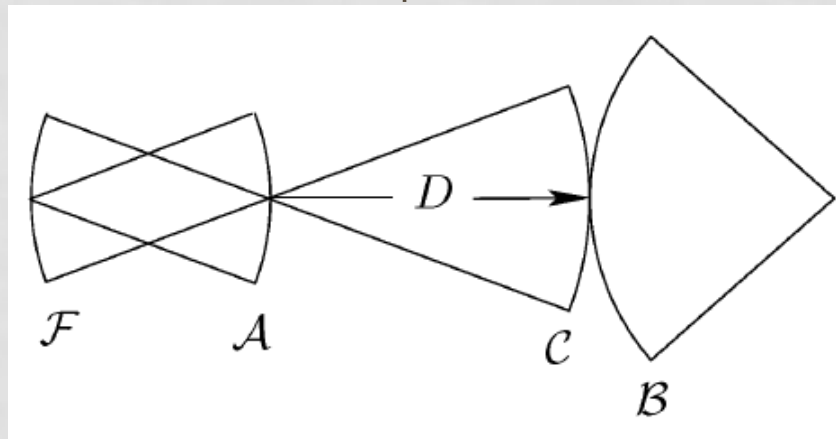
THE GENERAL TRANSFER



$$U_B(\mathbf{s}) = \frac{i}{\lambda D} \exp \left[-\frac{i\pi}{\lambda} \left(\frac{1}{R_B} + \frac{1}{D} \right) s^2 \right] \\ \times \int_{\mathbb{R}^2} \exp \left[-\frac{i\pi}{\lambda} \left(\frac{1}{D} - \frac{1}{R_A} \right) r^2 \right] \exp \left[\frac{2i\pi}{\lambda D} \mathbf{s} \cdot \mathbf{r} \right] U_A(\mathbf{r}) d\mathbf{r}$$

THE BONNET TRILOGY : FRESNEL FILTERING

- A Fresnel diffraction phenomenon from an emitter to a receiver can be composed into three operations
 1. An optical (spatial) Fourier transform (Fraunhofer diffraction)
 2. A spatial filtering, between Fourier sphere and a cardinal sphere tangent to the receiver
 3. A phase transformation (curvature transparency) between the cardinal sphere and the receiver



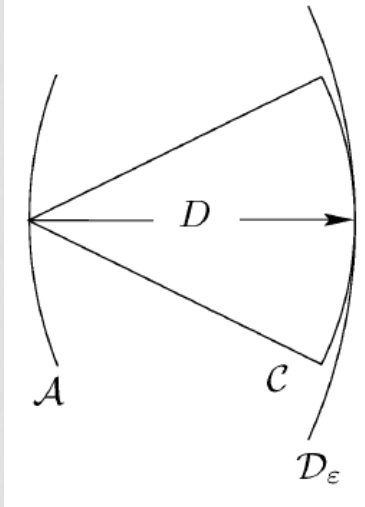
FRACTIONAL FOURIER:

GENERAL SPHERE TO SPHERE FIELD TRANSFER

- 2D Fractional Fourier Transformation Definition

$$\mathcal{F}_\alpha[f](\boldsymbol{\sigma}) = \frac{ie^{-i\alpha}}{\sin \alpha} \exp[-i\pi\sigma^2 \cot \alpha] \times \int_{\mathbb{R}^2} \exp[-i\pi\rho^2 \cot \alpha] \exp\left[\frac{2i\pi\boldsymbol{\sigma}\cdot\boldsymbol{\rho}}{\sin \alpha}\right] f(\boldsymbol{\rho}) d\boldsymbol{\rho}$$

- Metaxial interpretation $\varepsilon R_A > 0$ α in $]-\pi, \pi[$



$$\cot \alpha = \varepsilon \frac{1 - \mu}{\mu}, \quad \alpha D \geq 0$$

$$V_A(\boldsymbol{\rho}) = U_A(\sqrt{\lambda\varepsilon R_A} \boldsymbol{\rho}),$$

$$V_C(\boldsymbol{\sigma}) = U_C\left(\sqrt{\lambda\varepsilon R_A} \frac{\boldsymbol{\sigma}}{\cos \alpha + \varepsilon \sin \alpha}\right)$$

$$\boldsymbol{\rho} = \frac{1}{\sqrt{\lambda\varepsilon R_A}} \boldsymbol{r},$$

$$\boldsymbol{\sigma} = \frac{1}{\sqrt{\lambda\varepsilon R_A}} (\cos \alpha + \varepsilon \sin \alpha) \boldsymbol{s}$$

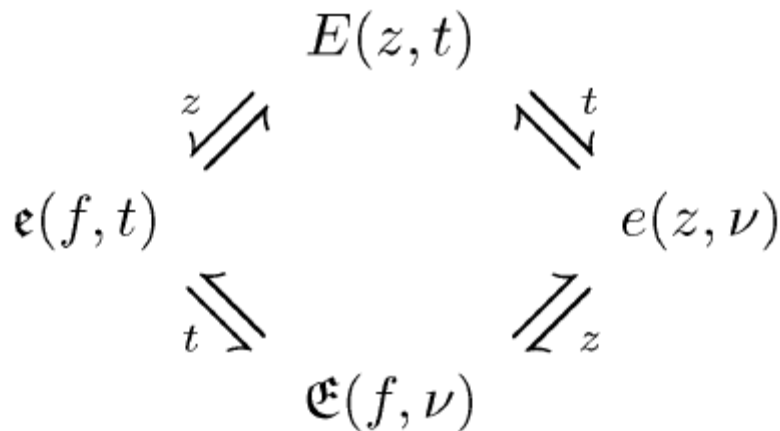
$$R_\varepsilon = \frac{\mu^2 + \varepsilon^2(1 - \mu)^2}{-\mu + \varepsilon^2(1 - \mu)} R_A = \frac{D^2 + \varepsilon^2(R_A - D)^2}{-D + \varepsilon^2(R_A - D)}$$

$$V_{D_\varepsilon}(\boldsymbol{\sigma}) = e^{i\alpha} (\cos \alpha + \varepsilon \sin \alpha) \mathcal{F}_\alpha[V_A](\boldsymbol{\sigma})$$

TIME DOMAIN FIELD TRANSFER
SECOND ORDER DISPERSIVE
OPTICAL FIBER CHANNEL

FIELD PROPAGATION ALONG THE LINE

- Field representation



$$e(f, t) = \int_{\mathbb{R}} E(z, t) e^{2i\pi f z} dz$$

$$E(z, t) = \int_{\mathbb{R}} e(f, t) e^{-2i\pi f z} df$$

$$e(z, \nu) = \int_{\mathbb{R}} E(z, t) e^{-2i\pi \nu t} dt$$

$$E(z, t) = \int_{\mathbb{R}} e(z, \nu) e^{2i\pi \nu t} d\nu$$

WAVE PACKET PROPAGATION

- We choose spatial (longitudinal) parameter $e(f, t) = a(f) e^{2i\pi\nu t}$

- Then

$$E(z, t) = \int_{\mathbb{R}} a(f) e^{2i\pi\nu t} e^{-2i\pi f z} df$$

- Because we choose spatial frequency as a parameter, then the (temporal) frequency is a function of f

$$E(z, t) = \int_{\mathbb{R}} a(f) e^{2i\pi[\nu(f)t - fz]} df$$

- We choose a wave packet function a defined by

$$a(\nu) = a[f(\nu)] \frac{df}{d\nu}$$

- Then

$$E(z, t) = \int_{\mathbb{R}} a(\nu) e^{2i\pi[\nu t - f(\nu)z]} d\nu$$

NARROW BAND PROPAGATION

- For narrow-band signal

$$\nu(f) = \tilde{\nu} + (f - \tilde{f}) \frac{d\nu}{df}(\tilde{f}) + \frac{1}{2}(f - \tilde{f})^2 \frac{d^2\nu}{df^2}(\tilde{f}) + \dots$$

- With the group velocity $V_g = \frac{d\nu}{df}(\tilde{f})$

- Narrow band propagation (the signal is delayed along the line)

$$\begin{aligned} E(z, t) &= e^{2i\pi(\tilde{\nu}t - \tilde{f}V_g t)} \int_{\mathbb{R}} \mathbf{a}(f) e^{2i\pi f(V_g t - z)} df \\ &= e^{2i\pi(\tilde{\nu}t - \tilde{f}V_g t)} E(z - V_g t, 0) . \end{aligned}$$

LINEAR FREQUENCY MODULATION

SECOND ORDER DISPERSIVE OPTICAL FIBER

- In linear frequency modulation, we have $\nu = \nu_0 + \nu_0 \theta t$

$$U(t) = e^{2i\pi\nu t} = e^{2i\pi\nu_0 t} e^{2i\pi\nu_0 \theta t^2} = U_0(t) e^{2i\pi\nu_0 \theta t^2}$$

$$e_z(\nu) = e_0(\nu) e^{-i\beta(\nu)z}$$

- Limited to second order, with $\beta_1 = (d\beta/d\nu)(\tilde{\nu})$ $\beta_2 = (d^2\beta/d\nu^2)(\tilde{\nu})$
- $\beta(\nu) = \beta_0 + \beta_1(\nu - \tilde{\nu}) + \frac{\beta_2}{2}(\nu - \tilde{\nu})^2$

$$\begin{aligned} E_z(t) &= e^{-i(\beta_0 - \beta_1 \tilde{\nu})z} \int_{\mathbb{R}} E_0(t') \sqrt{\frac{2\pi}{|\beta_2 z|}} e^{-is(\beta_2 z)\pi/4} \\ &\quad \times \exp\left[2i\pi\tilde{\nu}\left(t - t' - \frac{\beta_1}{2\pi}z\right)\right] \exp\left[\frac{2i\pi^2}{\beta_2 z}\left(t - t' - \frac{\beta_1}{2\pi}z\right)^2\right] dt' \\ &= e^{-i\beta_0 z} \int_{\mathbb{R}} E_0(t') \sqrt{\frac{2\pi}{|\beta_2 z|}} e^{-is(\beta_2 z)\pi/4} \exp[2i\pi\tilde{\nu}(t - t')] \\ &\quad \times \exp\left[\frac{2i\pi^2}{\beta_2 z}\left(t - t' - \frac{\beta_1}{2\pi}z\right)^2\right] dt'. \end{aligned}$$

The expression is very similar to Fresnel diffraction

COMPLEX ENVELOPE

$$A_z(t) = e^{-i\beta_0 z} \sqrt{\frac{2\pi}{|\beta_2 z|}} e^{-i\mathfrak{s}(\beta_2 z)\pi/4} \exp\left[\frac{2i\pi^2 t^2}{\beta_2 z}\right] \\ \times \int_{\mathbb{R}} A_0(t') \exp\left[\frac{2i\pi^2}{\beta_2 z} \left(t' + \frac{\beta_1}{2\pi} z\right)^2\right] \exp\left[-\frac{4i\pi^2 t}{\beta_2 z} \left(t' + \frac{\beta_1}{2\pi} z\right)\right] dt'.$$

- We define complex envelope function:

$$B_z(t) = \sqrt{\frac{2\pi}{|\beta_2 z|}} e^{-i\mathfrak{s}(\beta_2 z)\pi/4} \int_{\mathbb{R}} A_0(t') \exp\left[\frac{2i\pi^2}{\beta_2 z} (t - t')^2\right] dt'$$

$$B_z(t) = \sqrt{\frac{2\pi}{|\beta_2 z|}} e^{-i\mathfrak{s}(\beta_2 z)\pi/4} \exp\left[\frac{2i\pi^2 t^2}{\beta_2 z}\right] \\ \times \int_{\mathbb{R}} B_0(t') \exp\left[\frac{2i\pi^2 t'^2}{\beta_2 z}\right] \exp\left[-\frac{4i\pi^2 t t'}{\beta_2 z}\right] dt'$$

ENVELOPE COMPLEX PROPAGATION

- Define linear frequency modulation function

$$T(t) = e^{2i\pi\tilde{\nu}\theta t^2}$$

$$\Lambda = \frac{s(\beta_2)}{\tilde{\nu}} \quad R = \frac{s(\beta_2)}{2\theta} \quad \Lambda R = \frac{1}{2\tilde{\nu}\theta}$$

$$T(t) = \exp\left[\frac{i\pi}{\Lambda R}t^2\right]$$

- Complex envelope linearly modulated



$$Z = \frac{\beta_2 z}{2\pi\Lambda}$$

$$U_0(t) = B_0(t) \exp\left[\frac{i\pi t^2}{\Lambda R_0}\right]$$

$$U_z(t) = B_z(t) \exp\left[\frac{i\pi t^2}{\Lambda R_z}\right]$$

FRACTIONAL FOURIER TRANSFORM

REDUCED ENVELOPE COMPLEX PROPAGATION

- We define reduced envelope complex

$$V_0(\tau') = U_0(\sqrt{\varepsilon \Lambda R_0} \tau'),$$

$$V_z(\tau) = U_z \left(\frac{\sqrt{\varepsilon \Lambda R_0}}{\cos \alpha + \varepsilon \sin \alpha} \tau \right).$$

$$V_z(\tau) = e^{i\alpha/2} \sqrt{\cos \alpha + \varepsilon \sin \alpha} \exp \left[i\pi \left(\frac{1}{R_z} + \frac{1}{Z} \right) \frac{\varepsilon \Lambda R_0 \tau^2}{(\cos \alpha + \varepsilon \sin \alpha)^2} \right] \\ \times \int_{\mathbb{R}} \exp[i\pi \tau'^2 \cot \alpha] \exp \left[-\frac{2i\pi \tau \tau'}{\sin \alpha} \right] V_0(\tau') d\tau'.$$

- This is fractional Fourier transform if $\left(\frac{1}{R_z} + \frac{1}{Z} \right) \frac{\varepsilon \Lambda R_0}{(\cos \alpha + \varepsilon \sin \alpha)^2} = \cot \alpha$

$$V_z(\tau) = e^{i\alpha/2} \sqrt{\cos \alpha + \varepsilon \sin \alpha} \mathcal{F}_\alpha[V_0](\tau)$$

REFERENCES

1. Pierre Pellat-Finet, OPTIQUE DE FOURIER : Theorie Metaxial et Fractionnaire, Springer-Verlag France, 2009
2. Georges Bonnet, OPTIQUE METAXIAL (phenomenologie), polycopie, GESSY – Universite du Sud